# Semiconductor surface plasmons: a route to terahertz waveguides and sensors



Edmund Keith Stone School of Physics University of Exeter

A thesis submitted for the degree of *Doctor of Philosophy* March 2012

### Semiconductor surface plasmons: a route to terahertz waveguides and sensors

Submitted by Edmund Keith Stone to the University of Exeter as a thesis for the degree of Doctor of Philosophy in Physics 2012

This thesis is available for Library use on the understanding that it is copyright material and that no quotation from the thesis may be published without proper acknowledgement.

I certify that all material in this thesis which is not my own work has been identified and that no material has previously been submitted and approved for the award of a degree by this or any other University.

Edmund Keith Stone 2012

I would like to dedicate this thesis to Phyllis Mary Hamblen my loving gran who passed away during my second year.

#### Acknowledgements

During my time at Exeter many people have helped me, kept me sane and fed the insanity, all of which seem to be required for completing a PhD. To all those people, thank you. Lots are listed below, those who are not, I am sorry.

A special thank you must be said to my supervisor. Euan's enthusiasm for science is unshakable and affects all those around him. Even when the list of failed attempts appeared unending he was still positive. On a similar note, Professor Roy Sambles, thank you for suggesting my name to Euan and for your help in group meetings. Without my first summer project with you I'm not sure I would be writing a thesis currently. Tom Isaac, I am indebted to for teaching me all he knew about terahertz spectroscopy. Tom, the lock-in amplifier still only works if you catch it off guard when switching it on. During my time I have become more worried about the apparent sentience of some of our equipment.

I'd like to thank Professor Bill Barnes for many useful meetings whether it was 'back to basics,' terahertz or generation orientated. Always happy, and friendly with an insightful comment on the subject at hand. A great contributor to those early meetings was Baptiste, who along with Martin, Tim A, Tim T, John, Joe and Steph showed me just how interesting a PhD could be during my summer projects as an undergraduate. Others from the group who have moved on, Chris Burrows and James Parsons, you guys helped a lot with HFSS. James, your antics in the basement are still the topic of discussion, soon I feel the stuff of legend. Not forgetting Andy and George, your knowledge of the FIB and SEM was indispensable.

Matt L and Ian, as much as you might not like to admit it, you are both full of knowledge and happy to spread it about and help people, and between the unreproducible comments the discussions are nearly always useful. Two fellow terahertz spectroscopers, thanks for all your help Dmitry and Rostislav.

Matt B, Helen and Liz, since being undergrads together it's always been good to have a chat with you guys. Mel, thanks for helping me with maths when I hit a brick wall, even when the questions were ridiculously simple and I'd argued myself into a hole. It was always a pleasure helping you with your computer woes. Also, Mel, thanks for talking me through the panic, you'll be a great teacher! Celia, you should probably be mentioned several times in this section but tough. Thank you for all the talks, whether physics or unrelated, both at network and in uni you have always been a help and had sensible words of advice. Tim 'starkles' Starkey, you are a unique individual, whether you're fighting with HFSS or looking confused at experiments it's always interesting. I'd like to thank the other members of the electromagnetics group who I've not named specifically and also the guys in other groups who've been good to talk to especially graphene about fabrication issues. On that note Pete Hale and Sam Hornett, two more of Euan's students who have a wealth of knowledge and are happy to pass it on when it's needed. And of course the workshop guys, Nick Cole especially, with Pete Cann, and Adam and Dave for finding me Helium when I was desperate to get some experiments done before the laser moved to its new home.

I did some demonstrating whilst doing my PhD, second year labs would not have been the same if it were not for Martin Smith, Lee Summers, and Steve Hubbard (and others mentioned in other places here), thank you guys for keeping me sane, especially you Lee, if the astro doesn't work out you can always be a stand up.

There are several people who might be thinking they have been forgotten at this point, but a special mention has to go with those of us who spend our time in the basement of physics. Alfie, you are annoyingly funny, despite your similarities to Holmes. Steve, thanks for the discussions especially recently. Matt N, you missed the best times in the basement, sorry about that. Tom C, having you in the same office has been useful for someone to bounce ideas off, I hope I've helped you in the same way. Your proof reading of some of the chapters below was indispensable, thank you. Finally for physics, Caroline and Chris 'CJ' 'Holmesey' Holmes. I have spent most of the last 4 years with you guys. There is so much to say, the arguing, the word game, the list is pretty long. Thank you. Chris, you annoyed me greatly at some points, earthquakes, and using me as a drum kit are two things that spring to mind. But despite that you are a great friend and someone who I hope I do not lose contact with. Likewise Caroline, there are many things that could go here, none of which you would like me to write, I therefore will not. Just a heartfelt thank you, finish soon, we never did find the time to play a game of mini golf. I think that pretty much completes the list of people from physics, if I've missed you, sorry.

Outside physics there are a group of people who have helped keep me routed. Helped distract me when I've needed it and helped me get to this point in my life. My parents of course, they have always believed in me and supported me through everything I have done. If it was not for the way they have brought me up I do not know where I would be at this point. I should also thank my brother, Tris, for constantly suggesting I could wear things as a special hat. When Celia first met Tris, she was surprised as from when I talked about him she had assumed I was the older brother.

Some friends and influences from other places have also enabled me to get to this point. Friends from scouting, especially James and Matt, two of my oldest friends who are always good for a distraction. And Mr Bushrod, a great inspiration, my first real science teacher when I was around 9 years old, he inspired me to find out how the world worked with some truly great teaching.

Lastly, Sarah, thank you for being honest and for being there, when discussing me day with you, you always listened and then pointed out you didn't know what I was talking about. And of course thank you for making me some amazing food, your cheesecake is just fantastic.

#### Abstract

The terahertz regime has until recently been some what neglected due to the difficulty of generating and measuring terahertz radiation. Terahertz time domain spectroscopy has allowed for affordable and broadband probing of this frequency regime with phase sensitive measurements (chapter 3). This thesis aims to use this tool to add to the knowledge of the interactions between electromagnetic radiation and matter specifically in regard to plasmonics.

This thesis covers several distinct phenomena related to plasmonics at terahertz frequencies. The generation of terahertz radiation from metal nanoparticles is first described in chapter 4. It is shown that the field strength of the plasmon appears to relate to the strength of the generated field. It is also shown that the power dependence of the generated terahertz radiation is not consistent with the optical rectification description of this phenomenon. An alternative explanation is developed which appears more consistent with the observations. A simple model for the power dependence is derived and compared to the experimental results.

In chapter 5 the parameters that make good plasmonic materials are discussed. These parameters are used to assess the suitability of semiconductors for terahertz surface plasmon experiments. The Drude permittivity of InSb is measured here, leading to a discussion of terahertz particle plasmons in chapter 6. Finite element method modelling is used to show some merits of these over optical particle plasmons. This also includes a discussion of fabrication methods for arrays of these particles.

Finally, chapter 7 is a discussion of so called spoof surface plasmons. This includes some experimental work at microwave frequencies and an in depth analysis of open ended square hole arrays supported by model matching method modelling. Perfect endoscope effects are discussed and compared to superlensing. The thesis ends with a brief conclusions chapter where some of the ideas presented are brought together.

### Contents

C	Contents vi				
Li	List of Figures ix				
N	omer	nclatur	e x	viii	
1	Intr	oducti	ion	1	
<b>2</b>	The	eory		4	
	2.1	Surfac	e plasmons	4	
		2.1.1	Dispersion relation of surface plasmons	5	
		2.1.2	Surface plasmon polariton length scales	8	
		2.1.3	Surface plasmon coupling methods	11	
	2.2	Drude	e model	12	
	2.3	Locali	sed surface plasmons resonances	15	
		2.3.1	Quasi-static approximation	16	
3	Met	$\mathbf{thods}$		19	
	3.1	THz S	Spectroscopy	19	
		3.1.1	Generation and detection	23	
		3.1.2	Photodiode detectors	29	
	3.2	Sampl	le self adhesive support	32	
	3.3	Optica	al spectra	32	
	3.4	Finite	Element Method Modelling	33	
		3.4.1	Boundaries	<b>34</b>	
		3.4.2	Incident wave solutions	35	
		3.4.3	Eigenmode solutions	35	
		3.4.4	Element meshing	36	
		3.4.5	Uses	36	

### CONTENTS

4	Nai	oparticle plasmonic generation of terahertz radiation	37
	4.1	Samples	38
	4.2	Experimental set up	41
	4.3	Islandised films	42
	4.4	Nanosphere lithography arrays	43
		4.4.1 Thickness dependency	44
		4.4.1.1 Experimental results	44
		4.4.1.2 Numerical modelling	45
		4.4.2 Optical intensity dependency	49
		4.4.2.1 Simple model	49
		4.4.3 Angle dependence	53
	4.5	Conclusions and future proposals	55
<b>5</b>	Ter	hertz plasmonic materials	56
	5.1	Drude parameters of optical plasmonic metals	56
	5.2	Drude model and semiconductors	57
	5.3	Terahertz plasmonic materials	59
		5.3.1 Measuring the Drude parameters	59
		5.3.2 Results	61
		5.3.2.1 Temperature dependence	63
		5.3.3 Comparison between semiconductors	66
		$5.3.3.1$ Modelling of gratings with different semiconductors $\therefore$	<b>6</b> 8
		5.3.4 Photoexcitation of InSb	70
	5.4	Conclusions	76
6	Ter	hertz particle plasmons	77
	6.1	Fabrication	80
		6.1.1 Chemical wet etch	80
		6.1.2 Focused ion beam	82
		6.1.2.1 Post etching positioning	83
		6.1.2.2 Results	85
		6.1.3 Reactive ion etching and chromium masks	86
	6.2	Modelling results	89
		6.2.1 Particle pairs	91
		6.2.2 Varying sharpness	92
		6.2.3 Photomodulation	97
	6.3	Conclusions	98

7	Sur	face m	odes on open ended hole arrays	<b>99</b>
		7.0.1	Terahertz hole array measurement	99
		7.0.2	Motivations	100
	7.1	Theore	etical formalism	101
		7.1.1	Modal matching approach	101
		7.1.2	Dispersion relations	105
	7.2	Result	s and discussion $\ldots$	108
		7.2.1	Measurement of surface mode dispersion	108
		7.2.2	Asymptotic frequencies and mode splitting	112
	7.3	Role o	f surface modes in hole arrray transmission	114
		7.3.1	Near field transmission of hole arrays	114
		7.3.2	Far field transmission of hole arrays	117
	7.4	Conclu	usions	118
8	Con	clusio	ns and future work	<b>121</b>
	8.1	Public	eations	122
Re	efere	nces		123

## List of Figures

2.1	An incident p-polarised electromagnetic wave on the interface between	
	two materials of different permittivity ( $\epsilon_d$ and $\epsilon_m$ )	<b>5</b>
2.2	The dispersion relation for silver using a frequency dependent permit-	
	tivity found from the Drude model	8
2.3	The field confinement in the $z$ direction (normal to the surface) using	
	a permittivity found using the Drude parameters of silver, (a) into the	
	metal, and (b) into the dielectric, $\ldots$ $\ldots$ $\ldots$ $\ldots$ $\ldots$ $\ldots$	10
2.4	A schematic of the electric fields on the metal-dielectric interface for a	
	surface plasmon polariton, the decay into the metal and dielectric are	
	shown. For any moment in time the field must change direction along	
	the surface as the surface charge wave is longitudinal. $\ldots$ $\ldots$ $\ldots$	11
2.5	Schematic diagram of the blade coupling method, a form of scattering	
	coupling	12
2.6	The path traced by an electron in a metal according to the Drude model.	
	The electron undergoes many random scattering events resulting in a	
	random walk. The solid line indicates the situation where there is no	
	applied electric field and the dotted line the case when there is an applied	
	electric field resulting in a drift velocity $v_d$	14
2.7	Example complex permittivity for metals in the optical regime, the solid	
	lines represent the real part, and the dashed lines the imaginary part,	
	for (a) silver and (b) aluminium, found using the Drude model. $\ldots$	15
3.1	Schematic of the focusing terahertz time domain spectrometer (THz-	
	TDS) used in this project, the optical 800nm path is shown in red with	
	the terahertz beam path shown in green.	20

3.2	(a) Schematic of the collimated terahertz time domain spectrometer	
	(THz-TDS) in the transmission configuration used in this project, the	
	optical 800nm path is shown in red with the terahertz beam path shown	
	in green. (b) Schematic of the dry air box section of the collimated	
	THz-TDS in the reflection configuration.	20
3.3	Example traces from a THz-TDS setup, (a) shows the time domain trace,	
	whereas (b) shows the modulus of the complex frequency domain trace,	
	also shown is the phase measured through the system	23
3.4	An example photoconductive device, a bias voltage is applied across the	
	antenna gap, when the incident optical pulse excites charge carriers their	
	resultant acceleration results in an emitted pulse of terahertz radiation.	24
3.5	(a) The atomic structure of ZnTe . (b) The non-linear potential due	
	to the assymptric bond between the Zn and Te atoms. (c) the optical	
	excitation pulse, with the induced polarisation in a very thin non-linear	
	crystal, and the resulting terahertz pulse.	26
3.6	The dependence of the peak terahertz electric field on the fluence of	
	the generating optical pulse for ZnTe (solid line), a non-linear optical	
	rectification crystal, and a photoconductive antenna device (dotted line)	
	with a 4 kV bias voltage.	28
3.7	The three considered circuit diagrams for balanced photodiodes. It	
	should be noted that these diagrams are simplified containing only the	
	most important components. (a) The situation where the diodes are	
	unpowered and in parallel, their direct differential is then amplified. (b)	
	Has two 'pre-amps', one for each of the diodes, the differential opera-	
	tion is conducted using an op-amp. (c) The pre-amps and differential	
	amplifier are supplemented by a sample and hold stage to conduct the	
	stretching, this is then amplified by a fourth op-amp	31
3.8	The time domain transmission through 4 different self adhesive pres-	
	sure sensitive films, shown with a free space transmission through the	
	focussing THz spectrometer.	33
4.1	A schematic of the proposed generation method.	38
4.2	A scanning electron microscope image of an islandised film, with a mass-	
	thickness of 12 nm, the narrow channels the film are clearly illustrated.	39
4.3	Nanosphere lithography process for fabricating nanoscale triangles	40
4.4	A scanning electron microscope image of an array of triangles fabricated	
	using 780 nm diameter spheres using nanosphere lithography with a de-	
	posited metal depth of 50 nm	40

4.5	Optical extinction spectra for islandised films, and triangles fabricated	
	by nanosphere lithography.	41
4.6	The peak terahertz intensity measured from islandised films with differ-	
	ent mass-thickness.	42
4.7	(a) The time domain generated terahertz transmission from a 1 mm	
	$\langle 110\rangle$ ZnTe crystal (dashed line), and an array of triangular particles	
	fabricated from nanosphere lithography using 780 nm (solid line) diam-	
	eter spheres. (b) The frequency domain measurements of the generated	
	terahertz signal, also shown (as the dotted line) is an array of triangles	
	fabricated using and 390 nm diameter spheres.	43
4.8	Dependence of the peak generated terahertz intensity on the thickness	
	of particles at normal incidence (solid line, filled circles) and an inclina-	
	tion angle (angle of incidence) of $40^{\circ}$ (dashed line, filled squares). The	
	lines are shown to guide the eye. Both datasets are normalised to their	
	respective peaks and the normal incidence measurements are around a	
	factor of 10 smaller. The arrays was fabricated using 780 nm diameter	
	spheres	44
4.9	Characteristic field enhancement distributions for 40 nm thick curved	
	particles in a rhombic unit cell calculated using finite element modelling,	
	illuminated with a wavelength of 800 nm. (a) In cross-section, where	
	the inset shows the probe position for the maximum field enhancement	
	plots. (b) Looking in plan of the particles.	45
4.10	The dipolar plasmon resonance shown in the reflection spectra for dif-	
	ferent realistic modifications, modelled for triangles made from spheres	
	with a diameter of 780 nm.	46
4.11	The realistic modifications made to the model. (a) The starting point	
	with straight pointed triangles, (b) The difference between the straight	
	sided and sloped particles, (c) The modification to the straight triangles	
	to make the sides and points curved, and $(d)$ showing the modification	
	to the top (and bottom) edges to include a small amount of curvature.	47
4.12	Results of numerical modelling for curved triangular particles with rounded	
	edges, (a) reflection spectra for different thicknesses, $h = 20, 30, 40, 50, 60$ ,	
	and 90 nm, the arrow shows the direction of increasing thickness and $(b)$	
	the dependence of the peak field enhancements at 800 nm on the thick-	
	ness of particles at normal incidence.	48

4.13	Dependence of generated terahertz fluence as a function of the intensity	
	of the optical pulse incident onto the sample. ZnTe is shown as the circles	
	and dashed line, with the squares and solid lines representing measure-	
	ments of a NSL fabricated array, using spheres of diameter $780~\mathrm{nm}$ and	
	with a particle thickness of 50 nm	50
4.14	Measured angle dependence of the generated terahertz intensity, from	
	an array of 50 nm thick triangles fabricated using NSL with spheres of	
	diameter 780 nm. Also shown are two predictions from $(4.16)$ where	
	$n_4 = 1$ (the dashed line) and $n_4 = 5$ (the doted line)	54
5.1	Schematic of idealised conduction and valence electron bands for metals,	
	semiconductors and insulators. The Fermi level is indicated by $E_F$ , up	
	to which the electron states are filled. $\ldots$ $\ldots$ $\ldots$ $\ldots$ $\ldots$ $\ldots$	58
5.2	Schematic of the transmission through a slab of material with refractive	
	index $n_2$	60
5.3	(a) Real and (b) imaginary parts of the transmission of the two silicon	
	samples in the frequency domain. Also shown is the (c) real and (d)	
	imaginary parts of the permittivity found from the fitting parameters.	
	The experimental results are show as symbols and the lines are the fits.	
	Sample a is shown as the circles and solid lines and sample b as the	
	diamonds and dashed lines	62
5.4	(a) Time domain and (b) real and (c) imaginary parts of the Fourier	
	transformed frequency domain of the transmission through InSb sample	
	at room temperature. The time domain results are only experimental.	
	In the frequency domain the experimental results are represented as a	
	symbol for every $75^{\text{th}}$ data point and the lines are the fits. (d) Shows	
	the permittivity from the fitted Drude parameters, the solid line is the	
	real part and the dashed imaginary. The inset shows where the real	
	permittivity crosses -1 (shown as the dashed straight line)	64
5.5	The extracted (a) scattering rate and (b) plasma frequency from the	
	InSb sample at various temperatures	65

5.6	A log log plot of the surface plasmon frequency $(\omega_{sp})$ plotted against the	
	scattering rate $(\gamma)$ for different semiconductors found in the literature	
	(open shapes) and measured in this paper (filled shapes). The black data	
	points represent room temperature measurements, red ones are above	
	room temperature and blue below. It is discussed in the text that $\omega_{sp}$	
	needs to be in the range of typical terahertz time domain spectrometers	
	$(0.2 \rightarrow 2 \text{ THz}, \text{ the shaded region})$ with $\gamma$ as low as possible. It can	
	therefore be said that InSb is the best candidate for terahertz plasmonic	
	experiments. The bold silicon result is used below for modelling.	66
5.7	Schematic of the grating profile used for the finite difference time domain	
	modelling.	68
5.8	Finite difference time domain results for a monograting with a mark	
	to space ratio of 50%. The depth of the grating is varied showing the	
	changing coupling efficiency. The two panels are (a) InSb using the	
	Drude parameters found above and (b) silicon using Drude parameters	
	from the literature	69
59	Schematic of the focusing terahertz spectrometer with a third path al-	00
0.0	lowing the sample to be optically numbed	70
5 1(	) Time domain transmission measurement of InSb at 10 K with no pho-	
0.10	toexcitation (solid line) and illuminated by a 400 nm pulse (dashed line)	71
5 11	Decay of the induced reduction in the terahertz transmission with time	11
0.11	from an optical photoevcitation pulse. Each sub figure shows various	
	temperatures (indicated in the figure) for a single fluence. The fluences	
	are described in the text where (a) shows <i>iii</i> , and (b) shows <i>iu</i>	72
5 19	Three consecutive measurements of the photoevcitation decay using nump	12
0.12	$f_{\text{Honce}}$ is a temperature of 100 K. The first measurement is the solid	
	line with the first repeat being the dotted line and the second the deshed	
	and the second the dashed	79
5 14	One	15
0.13	Schematic of the conduction and valence bands showing fined electron	
	states as the system is photoexcited. (a) when the system starts without	
	being photoexcited, and (b) where the system has already been photoex-	70
		73
5.14	Decay of the induced reduction in the terahertz transmission with time	
	trom an optical photoexcitation pulse. Each sub figure shows various	
	fluences (indicated in each figure) for a single temperature. The temper-	
	atures are (a) 50K, and (b) 150K. $\ldots$	74

5.1	15 The first part of the decay trace measured at 10 K with pump fluence <i>iii</i> (symbols are the measured values). The rise time can be seen with a fitted Gaussian function (solid line) with a full width half maxima of	
	27 ps	75
5.1	16 Schematic showing impact ionisation. Here energy from the photoexcited electron $a$ is transferred to electron $b$ as it decays to $c$ , providing energy	
	for $b$ to be excited to $d$	75
6.1	1 A schematic of the unit cell for pillar localised surface plasmon resonances	
	with example field profiles.	78
6.2	2 Reflection spectra for InSb rods (a) connected by an InSb substrate with	
	varying particle heights, $h = 10 \ \mu\text{m}, 50 \ \mu\text{m}, 100 \ \mu\text{m}$ and 150 $\mu\text{m}$ , and	
	(b) disconnected square rods ( $h = 10 \ \mu m$ ). The dimensions of the rods	
	modelled are length 20 $\mu m,$ and width 10 $\mu m,$ with a separation between	
	the particles of 10 $\mu m.$ The array has a long axis pitch of 60 $\mu m$ and	
	short axis pitch of 20 $\mu$ m	79
6.3	3 Snapped side of InSb wafers at $35^{\circ}$ where the scale bars are 100 $\mu$ m. (a)	
	An unetched wafer, some stress damage from breaking the sample can	
	be seen, (b) a sample etched using a lactic and nitric acid mix, where	
	the etchant was only incident on the top surface, damage can be seen	
	over 200 $\mu$ m under the surface of the sample, and (c) the material has	
	been etched by $\approx 250 \ \mu m$ in nitric acid no damage other than that due	
	to the stress on breaking the sample can be seen	81
6.4	4 The surface of an InSb wafer after etching in nitric acid, (a) an optical	
	microscope image and (b) an SEM image, where the scale bar is 20 $\mu$ m.	82
6.5	5 The surface of an InSb wafer after etching in nitric acid with a chromium	
	mask of curved triangles being deposited before etching. This is an	
	optical microscope image, the base to point length of each triangle is	
	around 100 $\mu$ m	83
6.6	6 (a) Schematic view showing the four cuts required to cut a triangular	
	prism particle using a FIB on a flat InSb substrate, the cuts are tapered	
	slightly similar to the shape of the FIB. (b) Three remaining particles	
	on the side of an InSb wafer after teeth have been etched into the edge	
	of the angled side of the wafer before another cut is made to remove the	
	particles from the substrate. (c) An SEM image of a single particle cut	
	from the side of the wafer	84

6.7	An optical microscope image showing the positioning of the triangular	
	shaped particles cut using the FIB. The red mark is a positioning point.	
	The cuts into the substrate in the bottom right are where the particles	
	were made	85
6.8	(a) A SEM image of the array of triangles cut using a FIB and positioned	
	using a micropositioner. The outer ring is a metallic aperture. (b) The	
	normalised transmission intensity measured through the InSb particles.	86
6.9	Two regions of the same sample after the lift off procedure. Remaining	
	on the surface should be chromium triangles on a PMMA protective layer	
	on an InSb substrate. The difference between the regions is very visible.	
	The base to point length of the triangles is 100 $\mu$ m	88
6.10	Curling of the chromium layer on the photo resist is observed originating	
	at the edges of some triangles. The point of the stretched pipette is also	
	shown	88
6.11	The optical microscope setup used for scraping the chromium mask to	
	clean it. (a) An overview of the microscope, with the camera used shown.	
	(b) The hydraulic micropositioner that was utilised.	90
6.12	Part of a cleaned array of chromium triangles on a PMMA protective	
	layer on an InSb substrate.	91
6.13	Photographs of two InSb wafers with PMMA and chromium masks on	
	top. Both show burning from the RIE process. (a) shows the case with	
	the surrounding PMMA and (b) where the surrounding PMMA was first	
	removed	91
6.14	Optical microscope images of etched InSb pillars. (a) Side on at an angle	
	of $\approx 75^{\circ}$ and (b) top down.	92
6.15	Spectra of field enhancement in between the particles of particle pairs as	
	the separation distance is decreased. These are shown for (a) spherical	
	particles and (b) triangular particles.	93
6.16	Field plots of the time averaged electric field in a plane cutting through	
	the particles halfway from their base to their top. The spheres have a	
	diameter of 30 $\mu m$ and the triangles a base to tip length of 30 $\mu m$ with	
	a thickness of 10 $\mu m.$ $\ldots$ $\ldots$ $\ldots$ $\ldots$	93
6.17	Three different methods of altering the sharpness of particles, (a) trun-	
	cation, (b) rounding and (c) changing the internal angle $(\alpha)$ .	94
6.18	Rounded and pointed triangles overlaid showing the difference between	
	the two modelled particle shapes	94
6.19	Transmission response for rounded and pointed triangle arrays, the curved	
	triangles are shown as the solid line and the pointed as the dashed	95

6.20	Time averaged electric field profiles on a plane cut through the middle	
	of the particle, modelled in an infinite array. The particle thickness is	
	10 $\mu m,$ the long axis 61 $\mu m$ and they are modelled for an array with a	
	pitch of 150 $\mu m.$ These field plots are for pointed particles as shown, at	
	(a) 0.8 THz, (b) 1.025 THz, and (c) 1.175 THz	95
6.21	Time averaged electric field profiles on a plane cut through the middle	
	of the particle, modelled in an infinite array. The particle thickness is	
	10 $\mu m,$ the long axis 61 $\mu m$ and they are modelled for an array with a	
	pitch of 150 $\mu \mathrm{m}.$ These field plots are for curved particles as shown, at	
	(a) 0.85 THz, and (b) 1.125 THz	96
6.22	Optical microscopy images of two arrays of (a) triangles and (b) trun-	
	cated triangles fabricated using the same photolithography mask	96
6.23	(a) A schematic of the three triangle types reproduced on the photolithog-	
	raphy mask. (b) The transmission through modelled arrays of these three	
	triangles. The sharpest triangles ( $\alpha = 15^{\circ}$ ) are shown as the solid line,	
	the intermediary particles ( $\alpha = 30^{\circ}$ ) as the dashed line, and the broad	
	particles ( $\alpha = 60^{\circ}$ ) as the dotted line	96
6.24	Calculated transmission through arrays of InSb triangles, the dashed	
	line is the normal response with the solid line showing the effect of an	
	increase in carrier concentration of $1 \times 10^{16} \text{ cm}^{-3}$ an increase of 50%.	97
7.1	(a) Scanning electron microscope image of the array of dimples in alu-	
• • -	minium, with diameter $\equiv 50$ µm, depth $= 90$ µm, and pitch $= 90$ µm.	
	The sample array is square with the direction of propagation, shown by	
	the arrow, at 45° to the lattice vectors. (b) The measured (solid curved	
	line) and analytical modelled (dashed line) dispersion of the terahertz	
	hole array sample. (c) 3D schematic of the unit cell used in numerical	
	modelling	101
7.2	Schematic diagrams showing (a) the open hole array sample configura-	
	tion, (b) the experimental set up, and (c) the unit cell of the experimental	
	sample, with the propagation direction indicated	102
7.3	Measured dispersion relations for open and closed metallic hole arrays in	
	the microwave regime, with analytical and finite element method mod-	
	elling dispersions shown. The field profiles for the symmetric and anti-	
	symmetric surface modes are shown on open ended hole arrays, found	
	using the finite element method	110

7.4	The transmission measured across the surface of an array of open ended
	square holes defined by $d = 9.25$ mm, $a = 6.96$ mm, $h = 15$ mm, and
	$\epsilon_h = 2.29$ , normalised against the transmission across a flat metal surface. 111
7.5	Surface wave dispersion from the analytical model for open ended hole
	arrays, (a) where the thickness is varied, and (b) where the surrounding
	permittivity is asymmetric above and below the sample
7.6	Dispersion relations for an open ended hole array defined by $a = 6.96$ mm,
	$h = 15$ mm, and $\epsilon_h = 2.29$ when $d = 15$ mm (solid curved line) and
	d = 9.25mm (dashed curved lines)
7.7	(a) Spoof surface mode dispersion relations in the reduced zone repre-
	sentation. For hole arrays defined by $\omega_{co} < \omega_{diff}$ , a number of modes
	associated with each grating vector are observable in the $1^{st}$ Brillouin
	zone. For hole arrays with $\omega_{co} > \omega_{diff}$ (b), only the first two (surface
	plasmon) modes are defined below cut-off. (c) Transmission as a func-
	tion of hole height $h$ for an array with dimensions $a = 0.8d$ . One can
	clearly observe the multimodal transmission, with several peaks in the
	frequency region between the cut-off frequency (marked by a solid ar-
	row) and $\omega = \omega_{diff}$ . (d) Transmission as a function of hole height h for
	an array with dimensions $a = 0.4d$ , showing the transmission mediated
	by the symmetric and antisymmetric surface plasmon modes. The lower
	panels show transmission spectra at heights indicated by dotted arrows.
	All calculations are for $\epsilon_h = 1119$

## Definitions

FEM	Finite element method
FIB	Focussed ion beam
InSb	Indium antimonide
SEM	Scanning electron microscope
SPP	Surface plasmon polariton
TE	Transverse electric (s-polarised)
THz - TDS	Terahertz time domain spectroscopy
TM	Transverse magnetic (p-polarised)

Microwave regime	$1 \rightarrow 200 \text{ GHz}$
$Optical/Visible\ regime$	$300 \rightarrow 1000 ~\rm{THz}$
Terahertz regime	$0.2 \rightarrow 2 \text{ THz}$