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ABSTRACT

Density functional theory (DFT) study of polypyrrole-TiO₂ composites has been carried out to explore their optical, electronic and charge transfer properties for the development of an efficient photocatalyst. Titanium dioxide (Ti₁₆O₃₂) was interacted with a range of pyrrole (Py) oligomers to predict the optimum composition of nPy-TiO₂ composite with suitable band structure for efficient photocatalytic properties. The study has revealed that Py-Ti₁₆O₃₂ composites have narrow band gap and better visible light absorption capability compared to individual constituents. The simulated results of band structure (band gap, and band edge positions), molecular orbitals, and UV–vis spectra of the optimized nPy-Ti₁₆O₃₂ systems strongly support the existence of strong interactions between Py and TiO₂ in the composite. A red-shifting in λ_{\max} , narrowing band gap, and strong intermolecular interaction energy (–41 to –72 kcal/mol) of nPy-Ti₁₆O₃₂ composites confirm the existence of strong covalent type interactions. Electron–hole transferring phenomena are simulated with natural bonding orbital analysis where Py oligomers found as donor and Ti₁₆O₃₂ as an acceptor in nPy-Ti₁₆O₃₂ composites.

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1. Introduction

Nanostructured TiO₂ is one of the most investigated photoactive materials due to its excellent electronic structure and high stability to photocorrosion under redox conditions [1–3]. TiO₂ has drawn much attention due to its potential applications in solar energy conversion (such as photocatalytic water splitting and dye-sensitized solar cell), environmental cleaning (photocatalytic degradation of pollutants, self-cleaning, and water purification), and bio-sensing [4–7]. The solar energy conversion and photocatalytic efficiency of TiO₂ is limited due to its wide band gap (3.20 eV), as it can only absorb UV region of the solar spectrum [8]. Along with wide band gap, TiO₂ also suffer from high charge recombination rate (photogenerated electron–hole pair) which hindered the photocatalytic activity [9]. Strategies such as noble metal deposition as co-catalyst, metal or non-metal doping and photosensitization have been adopted to tailor band gap and enhance the charge transport properties to improve the photocatalytic performance of TiO₂ [10,11].

Conjugated organic polymer (COPs) [12–14] having spatially extended π -bonding system become potential energy materials due to their unique electrical and optical properties, such as high

photon absorption coefficients under visible light irradiation, high electron mobility, excellent stability, material diversity, mechanical flexibility, light weight, low-temperature processing, roll to roll printing and large-area capability [15]. They have been successfully applied in organic electronic devices and solar cells [16]. Conducting polymers are mostly *p*-type semiconductors which work as an electron donor in the *p*-*n* junction solar cell. Polyaniline (PANI), polypyrrole (PPy), polythiophene (PT), polyacetylene (PA), polythiophene (PT), polyparaphenylene (PPP), polyparaphenylen-enevylene (PPV), poly(3,4-ethylenedioxythiophene) (PEDOT), poly(*o*-phenylenediamine) (POPD) are prominent examples of COPs which are being explored in this area [12–14,17]. Recently, it has been found that composite of metal oxide and COP into intimate contact on both physical and electronic levels have significantly improved solar energy conversion and photocatalytic performance compared to metal oxide or COP individually [18]. Moreover, an increasing interest is found on the fabrication of COP-TiO₂ nanocomposites for photocatalytic degradation of environmental pollutants as well as water splitting to generate hydrogen [19,20].

Although, experimental study of PPy-TiO₂ as a photocatalyst is reported but facing lower photocatalytic efficiency which is due to lack of theoretical investigation [21–24]. The fully explored theoretical insight of optical, electronic structure, surface interactions, electroactivity and charge transfer mechanism between PPy and TiO₂ is essential to tailor and improve photocatalytic activity of PPy-TiO₂ nanocomposite. Our recent investigation on PPy and pyrrole oligomers (nPy, where *n* is a number of repeating units) has proved

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that nPy has excellent tunable, optical and electroactive properties [25–27]. In this work, we have detailed theoretical investigation of charge transfer mechanism between nPy and TiO₂ (molecular cluster; Ti₁₆O₃₂) in a connected junction such as nPy-TiO₂ nanocomposite. The present article is focused on the investigation of important photocatalytic parameters such as (i) interaction of Ti₁₆O₃₂ with PPy oligomers to find out electron–hole transformation, (ii) band structure alteration, (iii) mechanism of charge (electron/hole) transport (vi) optimisation of oligomeric length of Py for efficient performance and (v) structure-property relationship. The study is aimed to provide a better understanding of the electronic structure of nPy-Ti₁₆O₃₂ nanocomposite and will open new direction to synthesize an efficient nPy-TiO₂ nanocomposite photocatalyst with the optimized composition for visible light photocatalysis. Moreover, this work will also inspire computational scientists to explore structural property relationship of other COP-metal oxide nanocomposites.

2. Methodology

Interaction of molecular cluster of Ti₁₆O₃₂ with nPy oligomers (n = 3, 5, and 9) are performed with the help of density functional theory (DFT). All DFT calculations [28,29] were carried out on GAUSSIAN 09 [30] and the results were visualized through Gabedit [31], GaussSum [32], and GaussView [33]. DFT and time-dependent DFT (TD-DFT) calculations were performed at B3LYP with LanL2DZ basis set for the determination of electronic structure properties of nPy-Ti₁₆O₃₂, as an efficient photocatalyst. Oligomers up to eight repeating units can accurately represent its polymeric characteristics so, oligomeric size is restricted to nine repeating units [28,29]. A molecular cluster of Ti₁₆O₃₂ is added to the oligomeric backbone of nPy and then optimized. Geometries of nPy-Ti₁₆O₃₂ (n = 3 to 9) systems were optimized at B3LYP/LanL2DZ level and confirmed from frequency calculations. The non-bonding and bonding interaction energies in nPy-Ti₁₆O₃₂ are simulated with help of interaction energy (ΔE_{int}) and geometrical counterpoise correction (gCP). The gCP is able to treat both inter- and intra-molecular BSSE on the same footing at low computational cost. The B3LYP-gCP-D3 [34] interaction energy is simulated from the Grimme Web service [35]. The ΔE_{int} and gCP-D3 are calculated from equations 1 and 2, respectively.

$$\Delta E_{int} = E_{Ti_{16}O_{32}} + E_{nPy} - E_{nPy-Ti_{16}O_{32}} \quad (1)$$

$$\Delta E_{gCP-D3} = E_{(nPy-Ti_{16}O_{32})gCP-D3} - E_{(nPy)gCP-D3} - E_{(Ti_{16}O_{32})gCP-D3} \quad (2)$$

UV–vis spectra are simulated at hybrid TD-DFT with pseudopotential of LanL2DZ. Natural bonding orbitals (NBO) analysis performed for the charge simulations. All these calculations were performed in the gas phase.

Table 1Electronic properties of isolated TiO₂, Ti₁₆O₃₂ and Ti₂₈O₅₆ Anatase clusters.

Species	Band gap (eV)	HOMO (eV)	LUMO (eV)	Optical gap (nm)
Ti ₁₆ O ₃₂	3.06	−7.33	−4.33	333
Ti ₂₈ O ₅₆	2.61	−7.39	−4.78	471
Isolated TiO ₂	1.86	−5.28	−3.42	647

3. Results and discussion

3.1. Choice of TiO₂ cluster

To understand physical significance of nPy-TiO₂ at the molecular level, it is important to select an appropriate TiO₂ cluster as a representative of its bulk. The interactions between nPy and TiO₂ surface can be described in two different ways; the cluster model and periodic surface models [36]. For the surface representation of TiO₂, we have used cluster models, consisting of 48 and 84 atoms which are applicable to the molecular quantum chemistry [36]. Electronic properties such as HOMO, LUMO, band gap, and optical gap of isolated TiO₂, Ti₁₆O₃₂, and Ti₂₈O₅₆ clusters are simulated which are comparatively listed in Table 1. On the basis of band gap simulation and computational point of view, it is evident that Ti₁₆O₃₂ (band gap 3.06 eV) is a better representative of bulk anatase (Fig. 1). The current study is focused on Ti₁₆O₃₂ and their interaction with different Py oligomers.

3.2. Optimized geometries

Geometrical relaxation of nPy, Ti₁₆O₃₂, and nPy-Ti₁₆O₃₂ bound species are carried out at DFT- B3LYP/LanL2DZ level. Optimized geometric structure of 48 atoms (Ti₁₆O₃₂) is given in Fig. 1, while their comparative geometrical parameters are listed in Table 2. Analysis of the data (Table 2) proves that there is no significant difference between geometrical parameters of bulk (crystal/condensed phase) and molecular cluster (gas phase), confirming that Ti₁₆O₃₂ is the best representative of bulk TiO₂.

The optimized geometric structures of 7Py-Ti₁₆O₃₂ and 9Py-Ti₁₆O₃₂ along with their inter-molecular bond distances are shown in Figs. 2 and 3, respectively, while 3Py-Ti₁₆O₃₂ and 5Py-Ti₁₆O₃₂ are given in S1 and S2 of the Supplementary information. During the optimization, Ti₁₆O₃₂ circulated and finally attached at the middle of nPy oligomer by establishing strong non-covalent bonds with H and C atoms of Py. The interaction of Ti₁₆O₃₂ with nPy oligomers, results distortion in nPy geometries (Figs. 2 and 3), which is a direct consequence of their strong orbital overlapping. Six different types of inter-molecular interactions are observed in the case of the 7Py-Ti₁₆O₃₂ system; where Ti and O of Ti₁₆O₃₂ have made the strong electrostatic type of bonding with the C and H atoms of 7Py; ranging at ca. 1.90 to 2.80 Å. In all these composites, Ti atom of Ti₁₆O₃₂

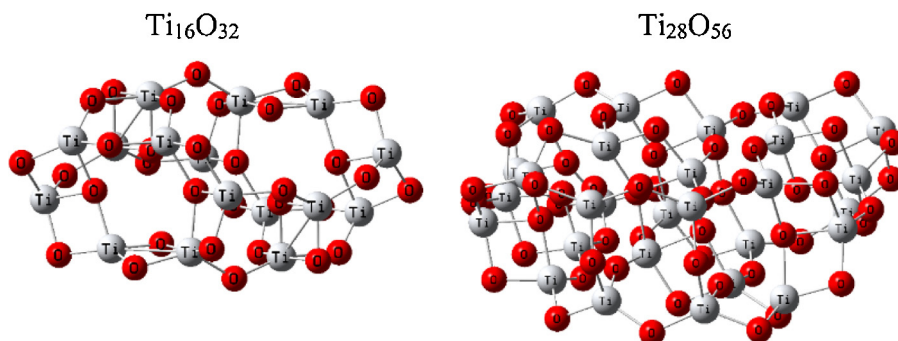
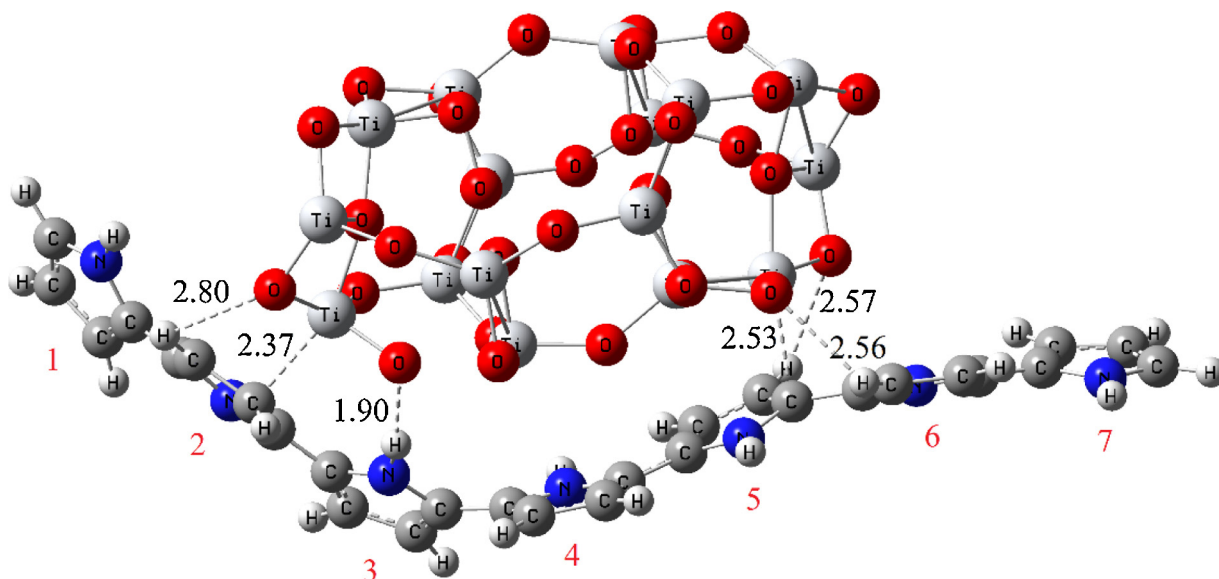
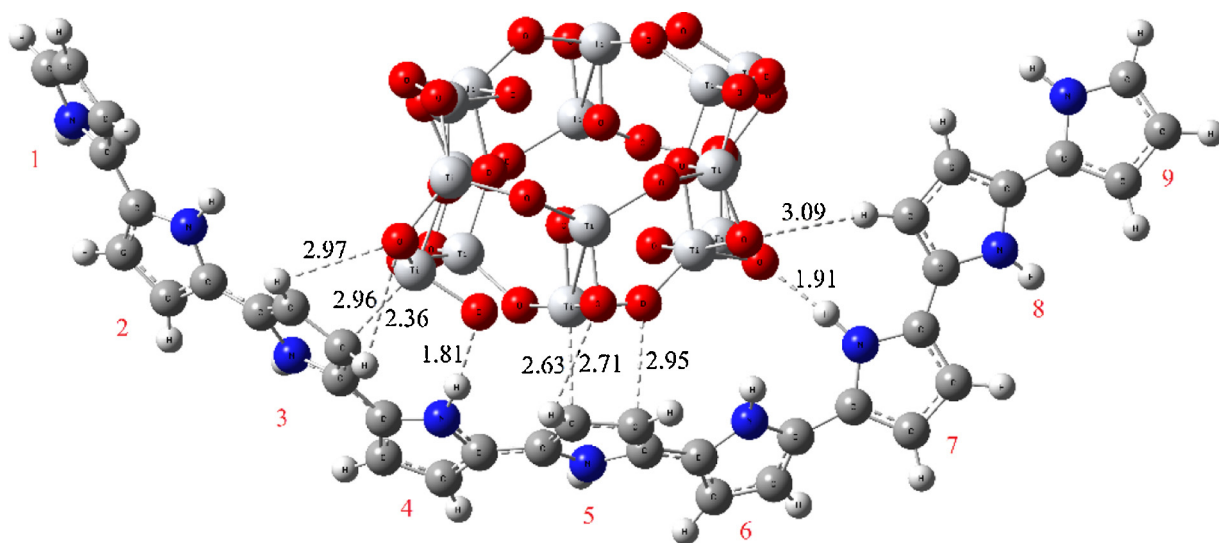
**Fig. 1.** Optimized Geometric Structure of Ti₁₆O₃₂ and Ti₂₈O₅₆ clusters.

Table 2
B3LYP/LanL2DZ ground states bond angles and distances in the Vacuum for Molecular TiO₂, cluster of Ti₁₆O₃₂, and in the Bulk for structural Rutile and Anatase.

Parameters	This study	TiO ₂ Vacuum	TiO ₂ Vacuum ¹	Rutile bulk properties ¹	Anatase bulk properties ¹
dTi–O	1.71/92 Å	1.63 Å	1.67 Å	1.96 Å	1.83 Å
dTi–Ti	2.78 Å	–	–	3.61 Å	3.03 Å
dO–O	2.46 Å	–	–	2.57 Å	2.29 Å
O–Ti–O	82.96°	110.7°	110.6°	81.6°	74.1°
Ti–O–Ti	102.82°	–	–	130.8°	105.9°

**Fig. 2.** Optimized Geometric Structure of 7Py-Ti₁₆O₃₂.**Fig. 3.** Optimized Geometric Structure of 9Py-Ti₁₆O₃₂.

makes an inter-molecular electrostatic bond with C of nPy and a Hydrogen bond between the O of Ti₁₆O₃₂ and H atoms of nPy.

The O–H inter-molecular bonds in the studied systems are strong bonds and can be regarded as covalent bonding. On the other hand, simulated Ti–C bond distances in the range of 2.30–2.70 Å confirm the formation of a strong electrostatic interaction. Jeffery et al. have reported that 2.2–2.5 Å of hydrogen bonding will be “strong, mostly covalent”, 2.5–3.2 Å as “moderate, mostly electrostatic” while 3.2–4.0 Å as “weak, mostly electrostatic” along with following bond energies, –40 to –14, –15 to –4, and <–4 kcal/mol, respectively [37]. Comparative analysis of

simulated inter-molecular bonding (including Fig. S1 and S2) led us to conclude that Ti₁₆O₃₂ make a strong composite with Py oligomers.

3.3. Interaction energy in nPy-Ti₁₆O₃₂ composites

Ti₁₆O₃₂ and PPy oligomers have a good interaction which is mostly covalent and electrostatic, simulated from interaction energy (ΔE_{int}) and geometrical counterpoise corrected ($\Delta E_{\text{gCP-D3}}$) methods (Table 3). The $\Delta E_{\text{gCP-D3}}$ method is employed to minimize the geometrical and dispersion factors. The result of this

Table 3
Inter-molecular interaction energy (ΔE_{int}), geometrical counterpoise corrected energy ($\Delta E_{\text{gCP-D3}}$), in kcal/mol, and NBO charges analysis of nPy-Ti₁₆O₃₂.

Species	ΔE_{int}	$\Delta E_{\text{gCP-D3}}$	Q_{NBO}
3Py-Ti ₁₆ O ₃₂	-42.54	-54.09	0.409
5Py-Ti ₁₆ O ₃₂	-34.22	-47.50	0.469
7Py-Ti ₁₆ O ₃₂	-41.10	-72.47	0.709
9Py-Ti ₁₆ O ₃₂	-41.47	-68.21	0.599

inter-molecular interaction energy proved that a very strong interaction is present in nPy and Ti₁₆O₃₂ species. In the 3Py-Ti₁₆O₃₂ composite, the ΔE_{int} is -42.54 kcal/mol while this interaction energy is simulated to be -54.09 kcal/mol with $\Delta E_{\text{gCP-D3}}$ method which is 11.55 kcal/mol higher than the ΔE_{int} method.

A comparatively weak inter-molecular forces are present in 5Py-Ti₁₆O₃₂ for which these energies are; -34.22 kcal/mol based on ΔE_{int} and -47.50 kcal/mol based on $\Delta E_{\text{gCP-D3}}$. This decrement in inter-molecular interaction energy is due to its geometrical shape compared to the 3Py-Ti₁₆O₃₂ composite. With chain length elongation, a slight difference in ΔE_{int} is observed as can be seen from data in Table 3. In the case of 7Py-Ti₁₆O₃₂, this inter-molecular non-bonding energy is about -41.10 kcal/mol, based on ΔE_{int} while -72.47 kcal/mol based on $\Delta E_{\text{gCP-D3}}$. This highest interaction energy ($\Delta E_{\text{gCP-D3}}$) in 7Py-Ti₁₆O₃₂ can be attributed to the suitable oligomeric length of PPy. Moreover, strong interaction of 7Py-Ti₁₆O₃₂ can be correlated to the fully relaxed geometric structure, which allows 7Py to wrap around Ti₁₆O₃₂ cluster much effectively compared to other compositions of nPy-Ti₁₆O₃₂ (Fig. 2). A similar but less pronounced trend is observed in 9Py-Ti₁₆O₃₂ composite, where the ΔE_{int} and $\Delta E_{\text{gCP-D3}}$ are -41.47 and -68.21 kcal/mol, respectively. Comparative analysis of the interaction energy of 7Py-Ti₁₆O₃₂ and 9Py-Ti₁₆O₃₂ indicate that 7Py make stronger nanocomposite. In summary, the amount of this interaction energy also evidences the existence of strong types of bonds between these two species which led to the confirmation of stable composite.

3.4. Natural bonding orbital analysis

Inter-molecular charge transfer between nPy and Ti₁₆O₃₂ are simulated with natural bonding orbital (Q_{NBO}) charge analysis at B3LYP with pseudopotential of LanL2DZ. As discussed elsewhere, these charge analyses are basis set dependent but with the same level of theory [such as B3LYP/6-31G (d), UB3LYP/6-31G (d), UB3LYP/6-311++G (d, p) or B3LYP/LanL2DZ] etc., for different structures the results would provide meaningful trends [38,39].

In connection with previous sections, Ti₁₆O₃₂ interacts with nPy oligomers through strong hydrogen (mostly covalent) and electrostatic (Ti–C). The analysis of NBO simulation indicates that Ti–C bond is the main charge transferring paths in all nPy-Ti₁₆O₃₂ composites. The net charge transfer in all nPy-Ti₁₆O₃₂ composites is listed in Table 3 while individual ring charges are listed in Table 4 along with Ti₁₆O₃₂ attached-Py rings. Analysis of data in Table 4 confirms the electron accepting nature of Ti₁₆O₃₂ in the composite as it withdraws electronic cloud density from Py oligomers to about 0.40 to 0.70 e⁻ which results in a cationic state in nPy oligomers, responsible for better electrical conductivity [40–42].

Table 4
NBO charge Analysis in unit of electron of nPy-Ti₁₆O₃₂ rings.

Species	Ring 1	Ring 2	Ring 3	Ring 4	Ring 5	Ring 6	Ring 7	Ring 8	Ring 9
3Py-Ti ₁₆ O ₃₂	0.228	0.11	0.071						
5Py-Ti ₁₆ O ₃₂	0.087	0.102	0.204	0.034	0.042				
7Py-Ti ₁₆ O ₃₂	0.065	0.16	0.103	0.108	0.121	0.105	0.047		
9Py-Ti ₁₆ O ₃₂	0.039	0.043	0.158	0.126	0.156	0.041	0.021	0.011	0.004

Table 5
Molecular Orbital Energy, Band Gap (in eV) and Dipole moment (Debye) of nPy and nPy-Ti₁₆O₃₂ composites.

Species	HOMO	LUMO	Dipole moment	Band Gap
Ti ₁₆ O ₃₂	-7.33	-4.27	0.0007	3.06
3Py	-4.71	-0.44	1.56	4.27
3Py-Ti ₁₆ O ₃₂	-5.67	-2.12	13.63	2.55
5Py	-4.45	-0.76	1.08	3.69
5Py-Ti ₁₆ O ₃₂	-5.04	-3.53	10.89	1.51
7Py	-4.34	-0.91	0.58	3.43
7Py-Ti ₁₆ O ₃₂	-5.24	-3.37	9.97	1.87
9Py	-4.28	-0.98	0.07	3.30
9Py-Ti ₁₆ O ₃₂	-4.42	-3.68	8.75	0.74

Finally, a unique charge distribution can be seen in the Py rings of 7Py-Ti₁₆O₃₂ system which suggest that 7Py-Ti₁₆O₃₂ is an optimum composition with greater charge transferring ability.

3.5. Electronic properties simulation

The electronic properties such as IP, EA, HOMO, LUMO, ESP, DOS and band gap of isolated as well as nPy-Ti₁₆O₃₂ composites was estimated at B3LYP/LanL2DZ level. The ΔSCF and negative of HOMO orbital energies are almost similar, especially for long chain systems. The IP energy is estimated from the negative of HOMO while the EA is obtained from negative of LUMO, using Koopman's theorem. Higher the EA and IP values of a chemical substance, greater will be its electroactivity and stability. The HOMO and LUMO energies (band edge positions) of isolated Ti₁₆O₃₂ are listed in Table 5 while frontier molecular orbitals are shown in Fig. 4.

Our simulated band gap value for Ti₁₆O₃₂ is 3.06 eV at B3LYP/LanL2DZ level, which is close to that of the observed data (3.20 eV) [40]. The difference is because of molecular cluster simulation as reported by Troisi et al. [41]; different simulation packages, cluster size, and level of theories give rise to different band gaps. Usually, the band gap decrease with increase in the size of the cluster and it is expected that this will also converge to the bulk value like that of the total energy [41]. Best reactive sites of both of the Ti₁₆O₃₂ and nPy oligomers are estimated from molecular electrostatic potential (MEP) as shown in Fig. 4 and 5. Based on these MEP plots, Ti₁₆O₃₂ and nPy oligomers interacted at a suitable distance.

3.6. Ionization potential

Interaction of Ti₁₆O₃₂ with nPy increases the IP of nPy oligomers indicating the increased stability of the resulted composite as can be seen from the delocalized contours (Fig. 5). Higher the IP more will be the stability of a chemical substance. In case of 3Py-Ti₁₆O₃₂ composite, the increase in IP is about 0.96 eV which consequently narrow its band gap from 4.27 to 2.12 eV (as discussed in band gap section below). A similar trend is observed in rest of the nPy oligomers in nPy-Ti₁₆O₃₂ composites, except 9Py-Ti₁₆O₃₂ composite, where negligible change (IP value) has been observed which can be attributed to lower interaction ability of their molecular orbitals. The increase in IP value of nPy-Ti₁₆O₃₂ composites is clearly depicted in Fig. 6.

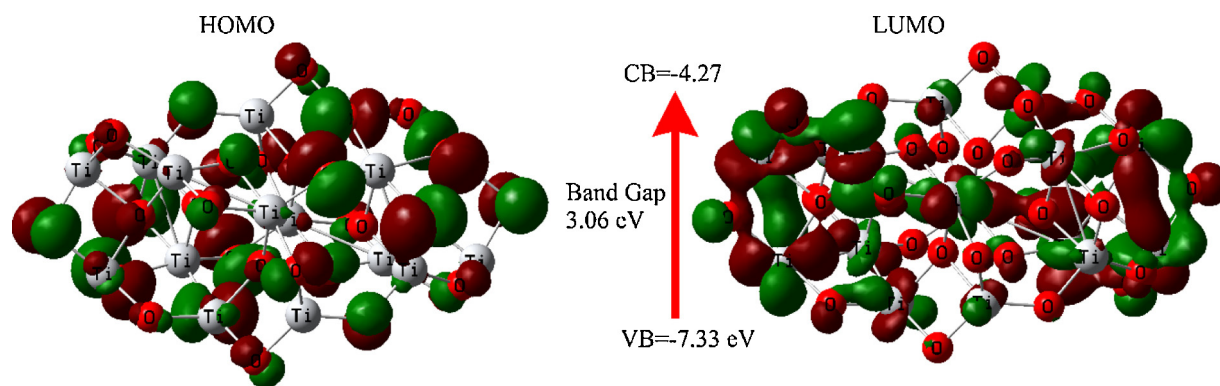


Fig. 4. Contours of the HOMO and LUMO of $\text{Ti}_{16}\text{O}_{32}$.

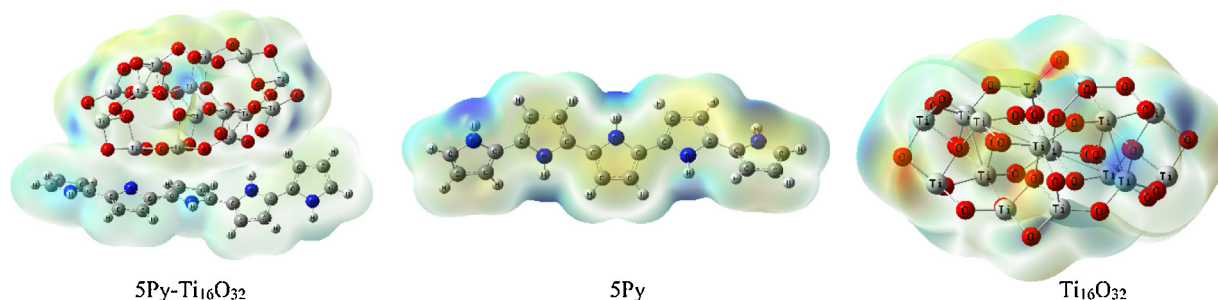


Fig. 5. Molecular Electrostatic Potential plots of $\text{Ti}_{16}\text{O}_{32}$, 5Py and $5\text{Py-Ti}_{16}\text{O}_{32}$.

3.6.1. Electron affinity

EA is estimated from the negative of LUMO, which found to be increasing in nPy oligomers upon interacting with $\text{Ti}_{16}\text{O}_{32}$. A substantial shift in LUMO to narrow the band gap is observed in nPy- $\text{Ti}_{16}\text{O}_{32}$ bounded species compared to individual $\text{Ti}_{16}\text{O}_{32}$ and nPy as shown in Fig. 7. $\text{Ti}_{16}\text{O}_{32}$ shifts the EA of 3Py from 0.44 to 2.12 eV, 2.77 eV in 5Py- $\text{Ti}_{16}\text{O}_{32}$, 2.46 eV in 7Py- $\text{Ti}_{16}\text{O}_{32}$ and 2.70 eV in 9Py- $\text{Ti}_{16}\text{O}_{32}$ composite. As discussed in the optimized geometric analysis, $\text{Ti}_{16}\text{O}_{32}$ has an excellent interaction with nPy oligomers which is also in agreement with the EA analysis. The overall increase in EA value proves that the nPy- $\text{Ti}_{16}\text{O}_{32}$ composites are cationic in nature where polaron and bipolaron states may exist.

3.6.2. Band gap and molecular orbitals energy

Both optical and electrical band gaps are simulated at ΔSCF (excitation energy with higher oscillator strength) and B3LYP (difference of HOMO-LUMO orbitals), listed in Tables 5 and 6. Analysis of the data (Table 5) clarifies that the band gap of nPy oligomers become narrow upon interaction with $\text{Ti}_{16}\text{O}_{32}$, which is an evidence of their excellent electroactive property in the resulted composite.

Frontier molecular orbitals and dipole moments of these interacting species before and after interaction are calculated at B3LYP/LanL2DZ level. The contours of HOMO and LUMO along with band gap of nPy and nPy- $\text{Ti}_{16}\text{O}_{32}$ composites are given in Figs. 8 and 9, respectively. It can be easily predicted from the data that molecular orbitals of $\text{Ti}_{16}\text{O}_{32}$ strongly interact with the

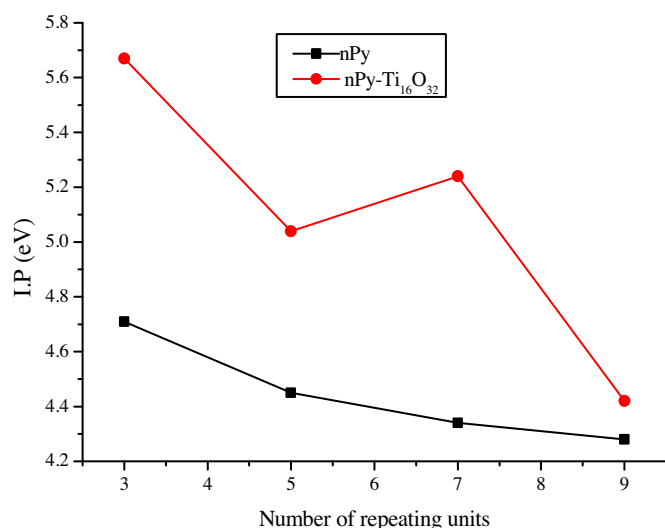


Fig. 6. Change in IP of nPy and nPy- $\text{Ti}_{16}\text{O}_{32}$ composites.

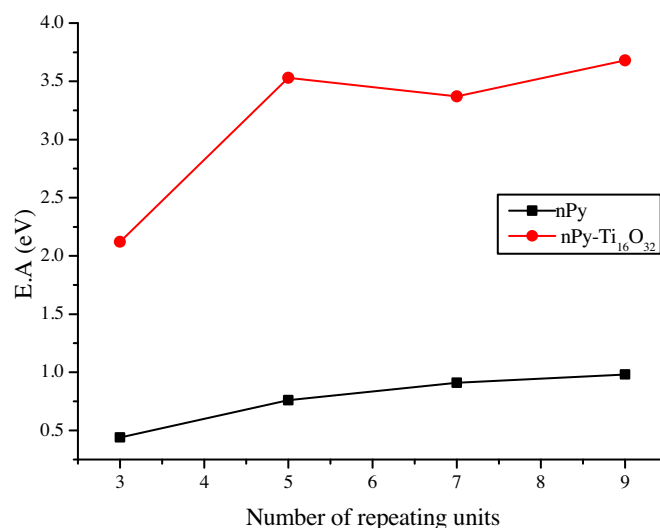


Fig. 7. Change in E.A. of nPy and nPy- $\text{Ti}_{16}\text{O}_{32}$ composites.

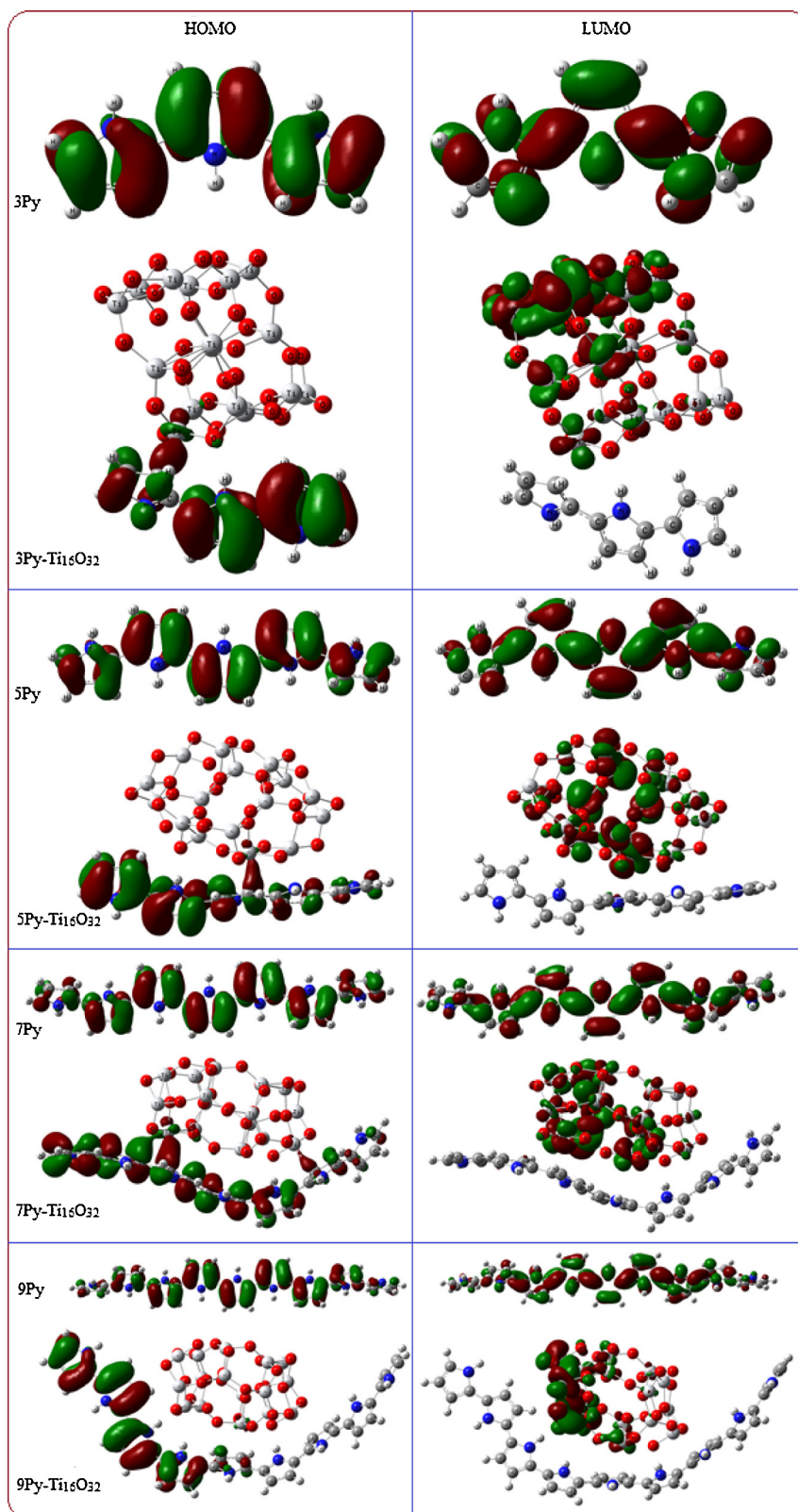


Fig. 8. Frontier molecular orbitals of $\text{Ti}_{16}\text{O}_{32}$, $n\text{Py}$ and $n\text{Py-Ti}_{16}\text{O}_{32}$ bounded complexes.

HOMO/LUMO of $n\text{Py}$ through establishing strong covalent types of bonding (Table 5). The strong covalent type of bonding indicate the stability of $n\text{Py-Ti}_{16}\text{O}_{32}$ composites which have an intermediate band gap and highest dipole moment.

The LUMO energy of $\text{Ti}_{16}\text{O}_{32}$ has considerably reduced from -0.44 to -2.12 eV while the band gap of Py reduced from 4.27 to 2.55 eV in all $n\text{Py-Ti}_{16}\text{O}_{32}$ composites as shown in Figs. 8 and 9. Dipole moment of $3\text{Py-Ti}_{16}\text{O}_{32}$ (13.63 Debye) clearly demonstrates a substantial increment in its electroactivity.

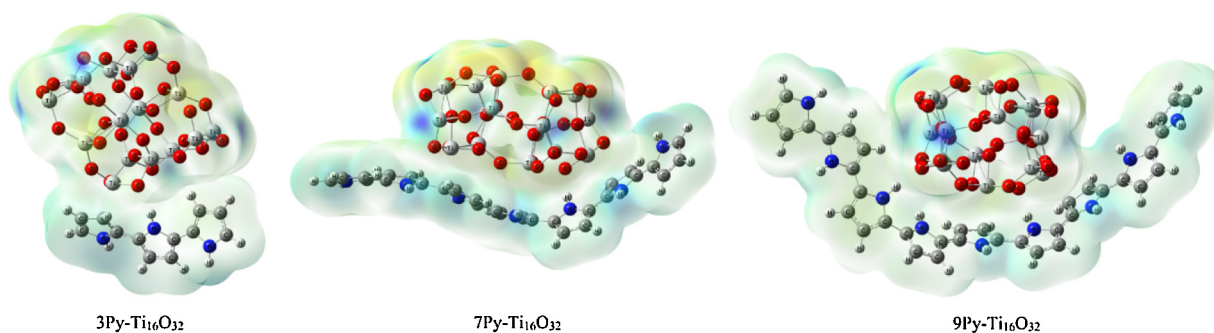


Fig. 9. Molecular electrostatic potential plot of 3, 7 and 9Py-Ti₁₆O₃₂ composites.

Moreover, narrowing band gap, have also a direct relationship with chain length elongation of Py as obvious from Fig. 10. Ti₁₆O₃₂ reduces the HOMO of 5Py from -4.45 to -5.04 eV (actually increase in negative charge), LUMO from -0.76 to -3.53 eV and band gap from 3.69 to 1.51 eV. Similar, but more pronounced reduction in these values can be seen in 7Py-Ti₁₆O₃₂, 0.90 eV in HOMO, 2.46 eV in LUMO, and a reduction of 1.56 eV in its band gap. Electronic cloud densities of the contours of HOMO and LUMO of 9Py-Ti₁₆O₃₂ compared to that of 9Py are enhanced to about 0.14 eV and 2.70 eV, respectively. While their dipole moment and band gap change from 0.07 to 8.75 Debye and 3.30 to 0.74 eV, respectively. The molecular electrostatic potential plot of 3, 7 and 9Py-Ti₁₆O₃₂ are shown in Fig. 9, which clearly visualize their electroactive nature.

Furthermore, the density of states plots for four these composites are correlatively shown in Fig. 10, which highlight a substantial variation in their band edges shifting compared to their individual constituents.

The homogeneous MEP plot (Fig. 9) of these composites confirm the excellent interaction between nPy and Ti₁₆O₃₂ cluster by mutually sharing of their electronic cloud densities. One can easily observe how the HOMO and LUMO of Ti₁₆O₃₂ (red spectra) and nPy (black spectra) interacted and resulted stable and ideal (align band edges with redox potential of water) nPy-Ti₁₆O₃₂ (blue spectra) composites. This theoretical investigation will guide and minimize the synthetic effort in term of interaction and mixing ratios. So, from experimental point view, we don't need to synthesize an

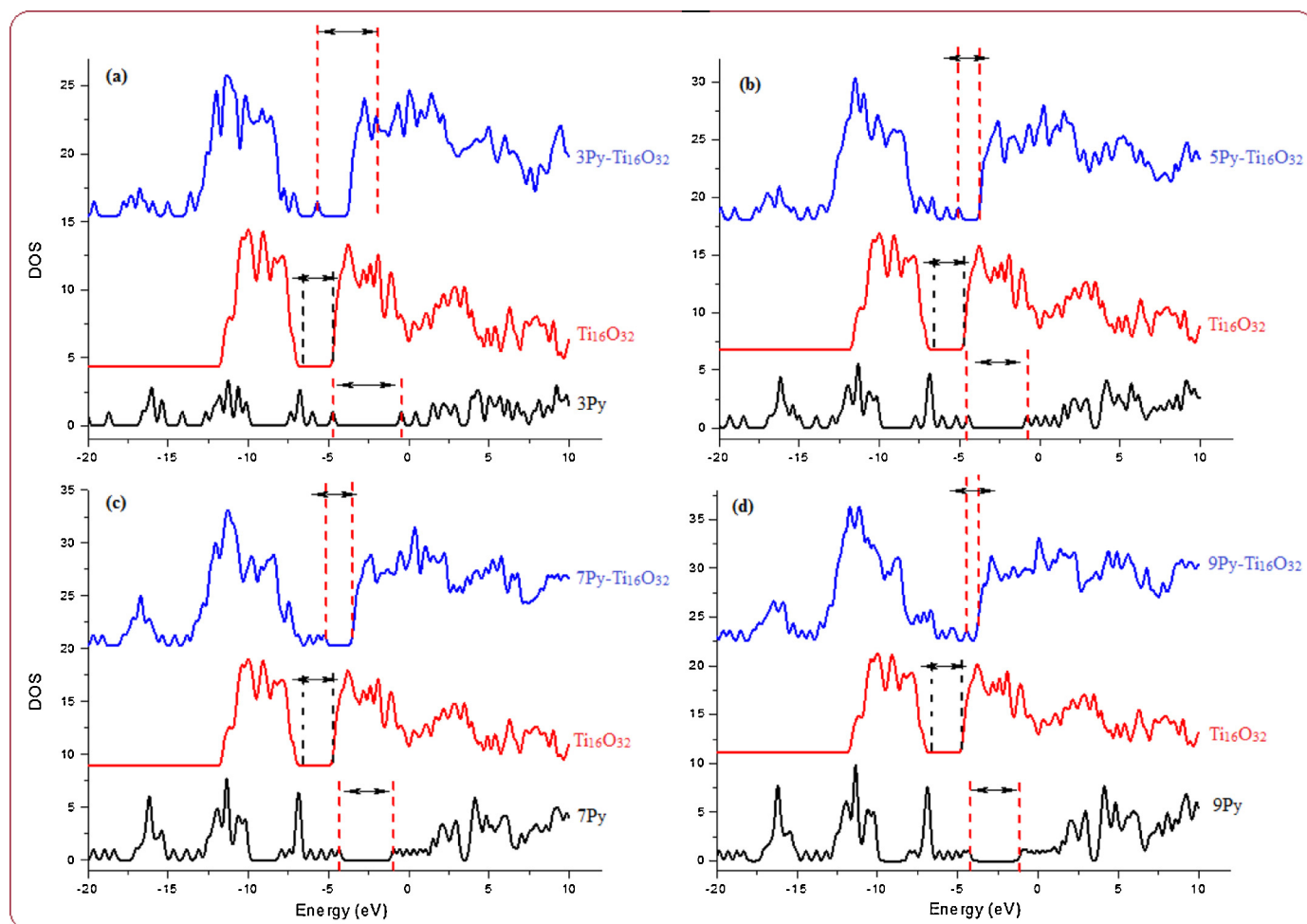


Fig. 10. Density of states plots of Ti₁₆O₃₂, nPy and nPy-Ti₁₆O₃₂ bounded complexes.

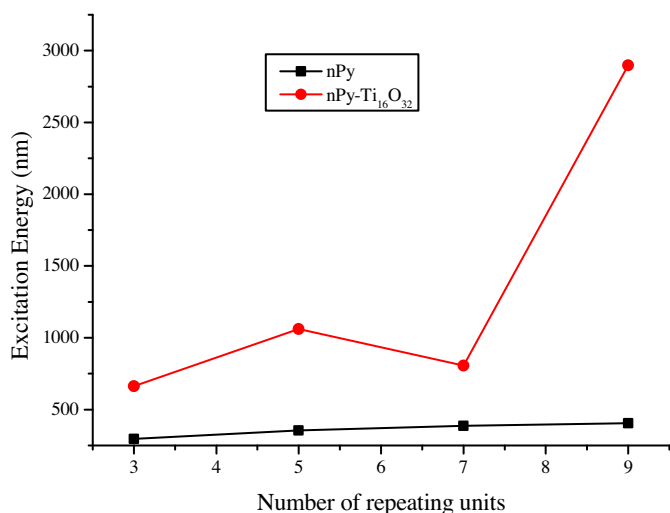


Fig. 11. First allowed electronic excitation energy of nPy and nPy-Ti₁₆O₃₂ composites.

infinite chain length of PPy, the routine oligomer (up to 6 or 8 repeating units) will be a best fit for the commercially available TiO₂ (anatase). Furthermore, the study also provides insight into improving the band structure, charge transport and determination of physical and chemical bonding between nPy and TiO₂ species.

3.7. UV-vis study

UV-vis spectra of nPy, Ti₁₆O₃₂, and nPy-Ti₁₆O₃₂ are simulated in the gas phase at TD-DFT/UB3LYP/LanL2DZ level. The first allowed electronic excitation energies are correlated with the experimentally observed λ_{\max} , which are listed in Fig. 11 and Table 6. Three prominent band peaks are found in the UV-vis spectra of nPy oligomers where the high wavelengths one is referred as λ_{\max} ; the transition of an electron from valence to conduction band. Interaction of Ti₁₆O₃₂ with nPy oligomers (Fig. 11 and Table 6) cause a red-shifting in the λ_{\max} of nPy. This red-shifting in λ_{\max} of all nPy oligomers illustrates the *n*-type doping nature of Ti₁₆O₃₂.

In the case of 3Py, the excitation energy of $\pi \rightarrow \pi^*$ transition increased from 296 to 663 nm in the resulted composite (3Py-Ti₁₆O₃₂), which evidences the establishment of strong bonding. As shown in Table 6, Ti₁₆O₃₂ increases the first allowed electronic excitation energy of 5Py to about 707 nm, 420 nm in 7Py, and 2493 nm in 9Py. This substantial increase can be attributed to improved conductivity/delocalization of the composites compared to isolated nPy oligomers (*vide infra*).

The individual Ti₁₆O₃₂ and nPy oligomers are unable to absorb in the visible range, however, nPy-Ti₁₆O₃₂ composites have excellent absorption capability in the visible region (Table 6) due to

Table 6

Calculated excitation energies, oscillator strengths, and molecular orbitals (MOs) of the first allowed singlet transition involved in the excitation for nPy and nPy-Ti₁₆O₃₂ composites.

Species	Energy (eV)	Wavelength (nm)	Oscillator Strength	MOs	Coefficient
Ti ₁₆ O ₃₂	3.71	333	0.01	S ₀ → S ₁	0.46
3Py	4.18	296	0.87	S ₀ → S ₁	0.70
3Py-Ti ₁₆ O ₃₂	1.86	663	0.005	S ₀ → S ₁	0.37
5Py	3.49	354	1.50	S ₀ → S ₁	0.70
5Py-Ti ₁₆ O ₃₂	1.16	1061	0.008	S ₀ → S ₁	0.50
7Py	3.19	387	2.08	S ₀ → S ₁	0.70
7Py-Ti ₁₆ O ₃₂	1.53	807	0.02	S ₀ → S ₁	0.68
9Py	3.05	405	2.67	S ₀ → S ₁	0.69
9Py-Ti ₁₆ O ₃₂	0.42	2898	0.0002	S ₀ → S ₁	0.70

a substantial decrease in the band gap. The red-shifting from the ultraviolet to visible elucidates/evidences the excellent photo-voltaic and photocatalytic activity of nPy-Ti₁₆O₃₂ composites over their individual constituents. Analysis of the simulation results led us to conclude that all nPy-Ti₁₆O₃₂ composites have a wide range of visible light absorption capability. Among the long chain oligomeric systems, the first allowed excitation energy of ($\pi \rightarrow \pi^*$ transition) of 9Py-Ti₁₆O₃₂ is more prominent compared to other compositions (Fig. 11) indicating that eight or nine repeating unit of PPy would be an excellent oligomeric chain length to develop an efficient visible light active photocatalyst (composite). This argument is also in good agreement with the other parameters discussed above.

4. Conclusions

Density functional theory study (DFT) of pyrrole-Ti₁₆O₃₂ bounded systems are carried out to find their interaction to tailor the best composite for the photodegradation of environmental pollutants and solar water splitting. Inter-molecular interaction energy in nPy-Ti₁₆O₃₂ composites is simulated in the range of -41 to -72 kcal/mol which confirmed the existence of strong covalent and electrostatic type bonding. This energy is simulated with the help of single point energy and $\Delta E_{\text{GCP-D3}}$ methods. So, after confirming the composite formation, band gap narrowing and better visible light absorption capability is observed compared to their individual nPy and Ti₁₆O₃₂ constituents. Electronic properties such as HOMO and LUMO of nPy, Ti₁₆O₃₂ and nPy-Ti₁₆O₃₂ composites estimated using B3LYP/LanL2DZ level, indicate excellent visible light absorption and charge transport efficiency of nPy-Ti₁₆O₃₂. Other electronic properties such as ESP, DOS, band gap, and UV-vis spectra also support the formation of efficient photoactive nPy-Ti₁₆O₃₂ composites. Moreover, UV-vis spectra of these composites predict the visible light absorption compared to nPy and Ti₁₆O₃₂. Electron-hole transferring define Py oligomers as a donor and Ti₁₆O₃₂ cluster as an acceptor in the resulted composites. Finally, the oligomeric length of eight/nine in the composite is found to be an optimum for the designing of an efficient photocatalyst. This theoretical investigation will minimize the synthetic effort such as mixing ratios of TiO₂ and PPy oligomers. Although, other components present in the cell might affect the nPy-TiO₂ interactions, but this basic theoretical study guides us about its possibility to be used as a photocatalyst. The study will also guide experimental scientists to improve the band structure, optical, physical and chemical properties of COP-TiO₂ and other metal oxide composites.

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Appendix A. Supplementary data

Supplementary data associated with this article can be found, in the online version, at <http://dx.doi.org/10.1016/j.snb.2016.10.019>.

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